

Published on Web 03/09/2009

Fluxionality of $[(Ph_3P)_3Rh(X)]$: The Extreme Case of $X = CF_3$

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Complexes of the type $[(R_3P)_3Rh(X)]$ are ubiquitous in organometallic chemistry and catalysis, the most renowned and widely used member of the family being Wilkinson's catalyst, [(Ph₃P)₃Rh(Cl)].¹ In the solid state, all of these species are squareplanar. While the vast majority retain the same rigid structure in solution, a small number of $[(R_3P)_3Rh(X)]$ complexes (X = H, Me, Ph) display equivalency of all three phosphine ligands.²⁻⁵ This unusual phenomenon has received surprisingly little attention in spite of its importance not only for basic science but also for the development of new applications in synthesis and catalysis. In one case, $[(Ph_3P)_3Rh(H)]$, the contrast between a square-planar geometry in the solid state and rapid intramolecular exchange in solution has been clearly established.⁴ The mechanism of this exchange, however, remains unknown, and it is still unclear how and why certain anionic ligands X promote the rearrangement whereas others confer stereochemical rigidity. Herein we report experimental and computational studies on a novel, uniquely fluxional complex, $[(Ph_3P)_3Rh(CF_3)]$, which maintains phosphine exchange at -100 °C. Our studies on this and related systems not only clarify the mechanism of the peculiar phosphine rearrangement in $[(R_3P)_3Rh(X)]$ but also provide new critical insights into the longpuzzling nature of bonding in perfluoroalkyl metal complexes.

Treating $[(Ph_3P)_3Rh(F)]^6$ with CF_3SiMe_3 in benzene resulted in a dark-red solution, from which uniformly shaped red-orange crystals were isolated upon addition of hexanes. X-ray analysis of the crystals (twice) revealed the structure of the product as *trans*-[(Ph_3P)_2Rh(CF_2)(F)] (1; Figure 1), a diffuorocarbene fluoride complex. The formation of **1** is likely mediated by $[(Ph_3P)_3Rh(CF_3)]$ (**2**), which undergoes α -F-elimination upon loss of one phosphine (eq 1).^{7–9}

While evidence was obtained for the same composition of **1** in the solid state in bulk,¹⁰ upon dissolution in benzene, toluene, or THF, preisolated **1** quickly equilibrated with a number of species, including **2**. A ¹⁹F and ³¹P variable-temperature (VT) NMR study indicated a complex system comprising several compounds in equilibrium, with some of the equilibria being fast and some slow on the NMR time scale at 25 °C. The presence of **2** pointed to phosphine dissociation from **1**. As **1** contains only two phosphines per Rh, we propose that isomers of dinuclear complexes [(Ph₃P)₂Rh₂(CF₂)₂(F)₂] could be produced.¹⁰ Remarkably, however, addition of excess PPh₃ to this multicomponent solution efficiently shifted all of the equilibria to **2** as the only NMR-detectable species.

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This allowed for high-yield (84%) isolation and full characterization of **2**, including X-ray analysis (Figure 1).



Figure 1. ORTEP drawings of 1 (left) and 2 (right). Selected bond distances (Å) and angles (deg) for 1: Rh–C, 1.820(3); Rh–F1, 1.994(2); Rh–P2, 2.335(1); Rh–P1, 2.349(1); C–Rh–F1, 177.0(1); P2–Rh–P1, 169.0(1); F3–C–F2, 100.0(2); F3–C–Rh, 130.8(2); F2–C–Rh, 129.1(2). For 2: Rh–C, 2.096(2); Rh–P2, 2.310(1); Rh–P1, 2.313(1); Rh–P3, 2.337(1); F1–C, 1.376(2); F2–C, 1.398(2); F3–C, 1.380(2); C–Rh–P1, 163.5(1); P2–Rh–P3, 156.5(1); F1–C–F3, 101.3(1); F1–C–F2, 101.7(1); F3–C–F2, 101.6(1).

Sharp doublets of quartets in the ¹⁹F and ³¹P NMR spectra¹¹ of **2** over the temperature range from 25 to -60 °C indicated fast *intramolecular*¹⁰ ligand exchange. Only at -100 °C (THF- d_8) was the exchange slow enough to observe two broad ³¹P doublets at 33.5 (2P, $J_{Rh-P} = 175$ Hz) and 30.5 ppm (1P, $J_{Rh-P} = 120$ Hz). Although temperatures below -100 °C were not attainable, and hence activation parameters could not be determined, magnetization-transfer experiments allowed for an exchange rate measurement of 12.1 s⁻¹ at -100 °C.

The methyl and phenyl analogues of 2, $[(Ph_3P)_3Rh(X)]$ (X = Me, Ph), were studied by VT ³¹P NMR for comparison. The reported^{6b} room temperature ³¹P NMR "unsymmetrical doublet" from the Ph species became symmetrical at 40 °C ($J_{P-Rh} = 163$ Hz), and below the coalescence point (around -20 °C), a complex second-order spectrum was observed, indicating apparent stereochemical rigidity.¹² For [(Ph₃P)₃Rh(Me)], the coalescence temperature was ~ 20 °C, and a well-resolved first-order A₂BX spectrum could be observed already at -10 °C. Measurements of the exchange rate over the temperature range from -10 to -50 °C allowed for the determination of activation parameters: $E_a = 16.5$ \pm 0.6 kcal mol⁻¹, $\Delta G^{\ddagger} = 12.9$ kcal mol⁻¹ (calculated at -30 °C), $\Delta H^{\ddagger} = 16.0 \pm 0.6 \text{ kcal mol}^{-1}$, and $\Delta S^{\ddagger} = 12.8 \pm 2.3 \text{ e.u. Using}$ this ΔS^{\ddagger} value with the exchange rate of 12.1 s⁻¹ measured for 2 at -100 °C (see above) led to an estimate of $\Delta H^{\ddagger} \approx 11.3$ kcal mol^{-1} for 2 under the careful assumption of similar entropies of activation for X = Me and CF_3 .

The above results point to the uniquely high fluxionality of **2** among known [(R₃P)₃Rh(X)] complexes. For R = Ph and X = Ph, Me, and H, the exchange rates are similar (e.g., 180 s⁻¹ at -10 °C for X = Me¹⁰ and 230 s⁻¹ at -13 °C for X = H⁴) and considerably lower than for X = CF₃ (12.1 s⁻¹ at -100 °C). To account for the

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facile exchange in **2**, we turned to density functional theory calculations¹³ to elucidate the mechanism of this process, and initial studies focused on the model system $[(H_3P)_3Rh(CF_3)]$ (**2**'). Computed bond lengths for **2**' agree well with those for **2**, although the distortion from square-planar geometry seen experimentally is not reproduced because of the use of PH₃ ligands (see below).



Figure 2. Computed reaction profile (kcal mol⁻¹) for intramolecular phosphine exchange in 2' with selected distances (Å) and angles (deg). Structures $2'_1/2'_2$ and $TS(2'-I')_1/TS(I'-2')_2$ are structurally equivalent, with the subscript indicating the P center trans to the CF₃ group. Energies shown in italics were computed at the MP2//BP86 level (see the text for details).

The reaction profile for phosphine exchange was investigated by reducing the trans C–Rh–P1 angle in $2'_1$ (i.e., with CF₃ initially trans to P1). This led to transition state $\mathbf{TS}(2'-\mathbf{I'})_1$ (E = +12.7 kcal mol⁻¹), in which CF₃ lies above the coordination plane (P1–Rh–C = 105°) and the cis phosphines move toward the vacant site (P2–Rh–P3 = 130°). $\mathbf{TS}(2'-\mathbf{I'})_1$ links to intermediate, $\mathbf{I'}$ (E =+11.5 kcal mol⁻¹) which resembles a trigonal bipyramid (TBP) with a vacant axial site trans to CF₃. The near- C_{3v} geometry of $\mathbf{I'}$ means that three equivalent TS structures can be accessed by increasing the relevant C–Rh–P angles: $\mathbf{TS}(2'-\mathbf{I'})_1$ returns CF₃ trans to P1, while $\mathbf{TS}(2'-\mathbf{I'})_2$ and $\mathbf{TS}(2'-\mathbf{I'})_3$ place CF₃ trans to P2 (see Figure 2) and P3, respectively. $\mathbf{TS}(2'-\mathbf{I'})_{1/2/3}$ are therefore high points on an energy surface that equilibrates all three phosphine ligands in 2' with (for the BP86 functional) $\Delta H_{calc}^{+} = 12.7$ kcal mol⁻¹.

Calculations on the full [(Ph₃P)₃Rh(CF₃)] system indicate a very similar topology for the phosphine-exchange surface. However, a considerable distortion away from square-planar is now computed,¹⁰ which is more in accord with the experimental data (Figure 1). The reactant is thus distorted toward the TS geometry, and a muchreduced barrier of only 4.0 kcal mol⁻¹ is calculated for the full system. The bulkier PPh₃ ligands therefore facilitate exchange, but ΔH^{\pm}_{calc} is significantly less than the approximate experimental value of 11.3 kcal mol⁻¹. This result was independent of the choice of functional and basis set, but an improved value of 8.7 kcal mol⁻¹ was obtained when the energies of the BP86-optimized species were recomputed at the MP2 level. The improved performance of the MP2//BP86 method was confirmed for [(Ph₃P)₃Rh(Me)] (ΔH^{\pm}_{calc} = 14.8 kcal mol⁻¹ vs 16.0 ± 0.6 kcal mol⁻¹ from experiment), so this approach was used subsequently.

Table 1 gives $\Delta H_{calc}^{\dagger}$ for a range of [(R₃P)₃Rh(X)] species, and these show good agreement with observed trends in fluxionality. In particular, the very high value of 48.5 kcal mol⁻¹ computed for $[(H_3P)_3Rh(Cl)]$ is consistent with the rigidity of Wilkinson's catalyst¹ at ambient temperature.¹⁴ Of the full systems, the TS for $[(Ph_3P)_3Rh(Cl)]$ unfortunately could not be located, but low barriers were found for $X = CF_3$, Me, Ph, and H, mirroring the fluxionality of these experimental systems. An intermediate barrier of 19.4 kcal mol⁻¹ was computed for X = CN, and experimentally, $[(Ph_3P)_3Rh(CN)]$ is rigid at room temperature.^{15,16}

Table 1.	$\Delta H_{calc}^{\ddagger}$ (MP2//BP86,	kcal mol ⁻¹	b for $[(R_3P)_3Rh(X)]$	
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Х	CF ₃	Me	Ph	Н	CN	CI
$\begin{array}{l} R = H \\ R = Ph \end{array}$	17.6 8.7	19.9 14.8	21.2 12.5	12.7 9.4	28.2 19.4	48.5 N/A

The steric factors that facilitate rearrangement in **2** are again evident for all of the $[(Ph_3P)_3Rh(X)]$ species in Table 1, with ΔH^{\pm}_{calc} always being less than those for the $[(H_3P)_3Rh(X)]$ congeners. Likewise, while $[(Ph_3P)_3Rh(Ph)]$ is fluxional at room temperature, sterically less encumbered $[(Me_3P)_3Rh(Ph)]^5$ and *cis*- $[(Ph_3P)_2(Ph_2PF)-$ Rh(Ph)] are not.^{6b,17} The size of X in $[(Ph_3P)_3Rh(X)]$ also plays an important role, with the reduction in barrier height upon introduction of the PPh₃ ligands being largest for X = CF₃ and Ph (~9 kcal mol⁻¹), less for Me (5.1 kcal mol⁻¹), and smallest for H (3.3 kcal mol⁻¹).

Steric effects alone, however, do not account for all of the observations: CF₃ and Br are approximately the same size, as are Me and Cl, and yet [(Ph₃P)₃Rh(Cl)] and [(Ph₃P)₃Rh(Br)] are rigid in solution¹ under conditions where $[(Ph_3P)_3Rh(Me)]$ and 2 are fluxional. Electronic factors are also important, and it is significant that the previously reported fluxional [(Ph₃P)₃Rh(X)] systems (X = Me, H, Ph) all feature ligands with strong trans influence.¹⁸ CF_3 is also considered to be a strong trans-influence ligand,¹⁹ and this is confirmed by the Rh-P1 distance of 2.31 Å in 2, which is as long as that measured in [(Ph₃P)₃Rh(H)].^{4,10} Figure 2 shows that as CF₃ moves into an axial position trans to a developing vacant site, the Rh-C bond shortens considerably, implying greater donation to the metal center. Similar changes are seen for all X, but one would expect the greater donor ability characteristic of ligands with the strongest trans influence to offer the greatest stabilization and thus produce lowest barriers. Weaker transinfluence ligands such as X = Cl are unable to stabilize the TS sufficiently, so intramolecular phosphine exchange by this mechanism becomes inaccessible.

Validation of these ideas comes from computed natural atomic charges for the reactant and TS structures (given in Table 2 for the simplified [(H₃P)₃Rh(X)] models). TS formation entails increased donation from the donor atom of X of $\sim 0.2e$. Despite this, the electron density on Rh actually diminishes in the TS, and it is the P centers that ultimately receive the extra charge density. This can be understood if the trigonal TS is considered to be a fragment of a TBP. Such a structure has two occupied d orbitals with significant σ -antibonding interactions with the phosphine ligands (consistent with increased Rh-P distances in both $TS(2'-I')_1$ and I' relative to 2', Figure 2). These d orbitals are approximately σ -nonbonding in the square-planar reactant, and thus TS formation entails a delocalization of electron density onto the P centers, each typically receiving $\sim 0.15e^{20}$ The loss of electron density at Rh can be mitigated by strong donation from X, and the highest charge density at Rh is found when $X = CF_3$, Me, H, and Ph, consistent with the strong trans influence of these ligands. The smaller charges on Rh in the TS for X = CN (despite its strong trans influence^{15,18}) and especially that for X = Cl correlate with the higher barriers computed in these cases. Judging only by the strongest negative charge on Rh (Table 2), one would expect the hydride [(Ph₃P)₃Rh(H)] to be the most fluxional species in the series. This

is not the case, however, possibly because of the steric effect (see above) that is counter-directing and particularly significant for the small H ligand. Entropic effects may also differ in this case.¹⁶

Table 2. Selected Computed Natural Atomic Charges (a) for [(H₃P)₃Rh(X)] Reactants and Transition States⁴

		reactants			transition states		
Х	q(Rh)	q(C∝/X)	q(Pav)	<i>q</i> (Rh)	q(C∝/X)	q(Pav)	
CF ₃	-0.52	+0.79	+0.23	-0.26	+0.93	+0.08	
Me	-0.48	-0.96	+0.23	-0.24	-0.73	+0.10	
Ph	-0.48	-0.22	+0.24	-0.20	+0.03	+0.08	
Η	-0.63	-0.09	+0.22	-0.45	+0.14	+0.10	
CN	-0.49	-0.02	+0.25	-0.17	+0.14	+0.10	
Cl	-0.45	-0.52	+0.25	-0.08	-0.33	+0.10	

^a Other centers display only minor changes in computed charge.¹⁰

Our computed data also allow us to address the factors underpinning the strong trans influence of the CF₃ ligand.^{19,21} Ligand trans influence is thought to be predominantly or exclusively controlled by field effects,²² so a strong trans influence implies strong electron donation. In complete contrast, the CF3 group is widely recognized in organic chemistry as a powerful electron acceptor.²³

The computed charges in Table 2 confirm that the Rh atom in 2' does indeed bear a large negative charge (-0.52). This actually exceeds that in the CH_3 analogue (-0.48), in spite of the opposite strong charges on the carbon atoms of the CF_3 (+0.79) and the CH_3 (-0.96) ligands (Figure 3). In fact, the charge distribution in 2' parallels the long-known²⁴ β -effect in fluorinated organic molecules: substitution of F for H on a carbon atom increases the negative charge on the next C (or H) atom. The explanation of the β -effect in terms of π -donation from the F atoms^{24a,c} accounts for the apparent flow of electrons toward the metal center in 2' (Figure 3). The charge data therefore shed light on the previously perplexing strong trans influence of CF₃ and other perfluoroalkyl ligands.^{19,21} Furthermore, an electrostatic attraction resulting from such $M^{\delta^--\delta^+}CF_3$ polarization could contribute to the M-C bond shortening in related polyfluoroalkyl complexes^{19,21} compared to their nonfluorinated counterparts.²⁵ The positive charge on the carbon atom, however, is probably stabilized by the same p_{π} backdonation mechanism from the fluorines.^{24,25b,c}



Figure 3. Natural atomic charges computed for $[(H_3P)_3Rh(CH_3)]$ and $2'^{10}$ and a resonance structure accounting for the β -effect.

In conclusion, the mechanism of the unusual intramolecular ligand exchange in complexes of the type $[(R_3P)_3Rh(X)]$ has been elucidated by experimental and computational methods. The rearrangement occurs via a distorted trigonal TS with the anionic ligand X in an axial position trans to a vacant site. Our results explain why certain $[(R_3P)_3Rh(X)]$ complexes (X = Cl,¹ Br, I, F,⁶ CN,¹⁵ OR,²⁶ NR₂²⁷) are stereochemi-cally rigid in solution, whereas others (X = Alk,^{2,5} Ar,^{2,6b} H^{3,4}) are fluxional under similar conditions. Exchange is governed by a combination of steric and electronic factors and is facilitated by bulkier ligands on the Rh as well as by strongly donating anionic ligands X that can stabilize the TS.²⁸ Our studies of the most fluxional new complex 2 ($X = CF_3$) have also clarified the previously puzzling strong trans influence and trans effect of R_f ligands as well as the nature of the bonding in perfluoroalkyl complexes of late transition metals.

Acknowledgment. This is DuPont CRD Contribution No. 8901. J.G. thanks Heriot-Watt University for support.

Supporting Information Available: Experimental and computational details and X-ray analysis data (CIF) for 1 (two polymorphs), 2, and [(Ph₃P)₂Rh(CO)(F)] and complete ref 13. This material is available free of charge via the Internet at http://pubs.acs.org.

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